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The influence of anesthesia and fluidstructure interaction on simulated shear stress patterns in the carotid bifurcation of mice

**Reference:**

De Wilde David, Trachet Bram, De Meyer Guido, Segers Patrick.- The influence of anesthesia and fluidstructure interaction on simulated shear stress patterns in the carotid bifurcation of mice

Journal of biomechanics - ISSN 0021-9290 - (2016), p. 1-7

Full text (Publishers DOI): <http://dx.doi.org/doi:10.1016/j.jbiomech.2016.06.010>

To cite this reference: <http://hdl.handle.net/10067/1354040151162165141>

1           **THE INFLUENCE OF ANESTHESIA AND FLUID-STRUCTURE**  
2           **INTERACTION ON SIMULATED SHEAR STRESS PATTERNS IN THE**  
3           **CAROTID BIFURCATION OF MICE**

4           **David De Wilde<sup>1</sup>, Bram Trachet<sup>1,2</sup>, Guido De Meyer<sup>3</sup>, Patrick Segers<sup>1</sup>**

5                   <sup>1</sup>IBiTech-bioMMeda, Ghent University-IMinds Medical IT, Ghent, Belgium

6                   <sup>2</sup>Department of Bioengineering, EPFL, Lausanne, Switzerland

7                   <sup>3</sup>Division of Physiopharmacology, University of Antwerp, Antwerp, Belgium

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9           **Keywords:** Carotid bifurcation, Mice, Wall shear stress, Computational Fluid Dynamics, Flu-  
10 id-Structure Interaction.

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12           Correspondence to:

13           Bram Trachet

14           bioMMeda, De Pintelaan185-blokB5, B-9000 Gent, Belgium

15           Tel: + 32 (0)9 332 46 21

16           Fax: +32(0) 93324159

17           Email: Bram.Trachet@ugent.be

18  
19           Word Count: 3427

22 **ABSTRACT**

23

24 **Background:** Low and oscillatory wall shear stresses (WSS) near aortic bifurcations have  
25 been linked to the onset of atherosclerosis. In previous work, we calculated detailed WSS pat-  
26 terns in the carotid bifurcation of mice using a Fluid-Structure Interaction (FSI) approach.  
27 We subsequently fed the animals a high-fat diet and linked the results of the FSI simulations  
28 to those of atherosclerotic plaque location on a within-subject basis. However, these simula-  
29 tions were based on boundary conditions measured under anesthesia, while active mice might  
30 experience different hemodynamics. Moreover, the FSI technique for mouse-specific simula-  
31 tions is both time- and labour-intensive, and might be replaced by simpler and easier Compu-  
32 tational Fluid Dynamics (CFD) simulations. The goal of the current work was (i) to compare  
33 WSS patterns based on anesthesia conditions to those representing active resting and exercis-  
34 ing conditions; and (ii) to compare WSS patterns based on FSI simulations to those based on  
35 steady-state and transient CFD simulations. **Methods:** For each of the 3 computational tech-  
36 niques (steady state CFD, transient CFD, FSI) we performed 5 simulations: 1 for anesthesia, 2  
37 for conscious resting conditions and 2 more for conscious active conditions. The inflow, pres-  
38 sure and heart rate were scaled according to representative in vivo measurements obtained  
39 from literature. **Results:** When normalized by the maximal shear stress value, shear stress pat-  
40 terns were similar for the 3 computational techniques. For all activity levels, steady state  
41 CFD led to an overestimation of WSS values, while FSI simulations yielded a clear increase  
42 in WSS reversal at the outer side of the sinus of the external carotid artery that was not visible  
43 in transient CFD-simulations. Furthermore, the FSI simulations in the highest locomotor ac-  
44 tivity state showed a flow recirculation zone in the external carotid artery that was not visible  
45 under anesthesia. This recirculation went hand in hand with locally increased WSS reversal.  
46 **Conclusion(s):** Our data show that FSI simulations are not necessary to obtain normalized

47 WSS patterns, but indispensable to assess the oscillatory behavior of the WSS in mice. Flow  
48 recirculation and WSS reversal at the external carotid artery may occur during high locomotor  
49 activity while they are not present under anesthesia. These phenomena might thus influence  
50 plaque formation to a larger extent than what was previously assumed.

51

## 52 INTRODUCTION

53 The hypothesis that low and oscillatory wall shear stress (WSS) has an atheroprone effect on  
54 the endothelial cells lining the inner arterial wall is generally accepted (Asakura and Karino,  
55 1990; Chatzizisis et al., 2007; Ku et al., 1985; Peiffer et al., 2013b). In mouse models of ath-  
56 erosclerosis, it has been shown that atherosclerotic lesions preferentially develop at branching  
57 points, bifurcations and at the bends of arteries (Nakashima et al., 1994). To be able to deter-  
58 mine murine WSS directly, highly accurate and therefore challenging measurements of the  
59 blood velocity near the arterial wall would be needed. A commonly used alternative is to use  
60 computer simulations to derive the WSS from the spatial velocity gradient perpendicular to  
61 the arterial wall. The simplest type of numerical simulations to calculate WSS in arteries are  
62 steady state CFD calculations. In these calculations either the peak or mean velocity (or flow)  
63 is applied as in-and outflow at the boundaries and the time-related variability is ignored.  
64 Transient CFD simulations also take into account the time dependency of the blood flow.  
65 Typically the velocity (or flows) are measured over an entire heart cycle and imposed as  
66 boundary conditions. In this case the time-averaged WSS (TAWSS) is used to quantify the  
67 shear stress patterns. Transient CFD allows to assess the oscillatory behavior of the flow and  
68 the WSS, which is assumed to play an important role in the atheroprone effect of low WSS on  
69 the endothelium. An important parameter in this respect is the Oscillatory Shear Index (OSI).  
70 A final step is the inclusion of the mechanical behavior of the arterial wall into the model, re-  
71 sulting in FSI simulations that take both the hemodynamics and the buffering effect of the ar-  
72 teries into account. This adds an extra layer of complexity to the modelling methodology, but  
73 also results in more realistic simulations.

74 Computer simulations of the mouse vasculature have already been performed by many au-  
75 thors. Most of these publications focused on the methodology and the feasibility to calculate

76 WSS in mice and did not link the found patterns (quantitatively or qualitatively) to biological  
77 responses or processes (Feintuch et al., 2007; Greve et al., 2006; Huo et al., 2008; Trachet et  
78 al., 2011a; Trachet et al., 2009; Van Doormaal et al., 2014). Nevertheless steady state CFD  
79 WSS patterns in the common carotid artery have been linked to the uptake of plasma macro-  
80 molecules (Mohri et al., 2014) or local wall thickness (Cheng et al., 2004). Transient CFD  
81 simulations have been used to link TAWSS and OSI patterns to aneurysm progression on a  
82 within-subject basis (Ford et al., 2011; Trachet et al., 2011b). In the carotid artery, transient  
83 CFD has been used to link wall shear stress patterns with plaque development (Assemat et al.,  
84 2014). In order to get access to the pre-disease geometry, the authors virtually removed the  
85 plaque from the model. Other studies linked wall shear metrics from transient CFD with  
86 plaque progression markers such as presence of fat in the wall (Suo et al., 2007) or the expres-  
87 sion of VCAM-1 (Hoi et al., 2011). In these studies, the CFD simulations were performed in  
88 another group of animals than the group that was used to determine the progression marker(s).  
89 Recently we have published a study presenting the first mouse-specific FSI simulations in the  
90 murine carotid bifurcation (De Wilde et al., 2015a), and we subsequently linked the shear  
91 stress patterns to plaque formation on a within-subject basis (De Wilde et al., 2015b).

92 When interpreting data from computer simulations in mice it is important to keep in mind that  
93 the used computational technique may influence the outcome of the simulations. When com-  
94 paring steady state and transient CFD in the mouse aortic arch only a small difference in abso-  
95 lute value of TAWSS was found (Feintuch et al., 2007). However, the effect of FSI on wall  
96 shear stress computations in mice has, to the best of our knowledge, never been investigated.  
97 Moreover, the boundary conditions used for these simulations are based on measurements ob-  
98 tained while the mice were under anesthesia. Anesthesia lowers the mean arterial pressure  
99 (MAP), the cardiac output and the heart rate (Janssen et al., 2004; Komarek, 2007). When the  
100 activity level of a mouse increases, these parameters will increase correspondingly. This can

101 considerably influence the simulated flow patterns (Mohri et al., 2014). Intrigued by these re-  
102 sults, we performed a parameter study in which we investigated not only the influence of dif-  
103 ferent simulation techniques (steady state CFD, transient CFD, FSI) but also the influence of  
104 anesthesia and locomotor activity on simulated wall shear stress patterns in the carotid bifur-  
105 cation of the mouse.

106

## 107 MATERIAL AND METHODS

### 108 Experimental protocol

109 The study was based on data acquired in a female ApoE<sup>-/-</sup> mouse that was fed a Western type  
110 diet (TD88137, Harlan Teklad, USA) ad libitum from the age of 6 weeks until 16 weeks. The  
111 housing of the animal and all the experiments were the same as in our previous publications  
112 (De Wilde et al., 2015a; De Wilde et al., 2015b). All experiments were approved by the ani-  
113 mal ethics committee of Ghent University and conducted according to their guidelines. At  
114 week 16 the mouse underwent an imaging protocol. Using a high-frequency ultrasound scan-  
115 ner (Vevo 2100, Visual Sonics, Canada) with a linear array probe (MS550D, 22-55 MHz), the  
116 blood flow velocity was assessed in the common, external and internal carotid arteries of the  
117 left carotid bifurcation with Pulsed Wave Doppler. The distension of the artery was measured  
118 using M-mode. A contrast-enhanced (100  $\mu$ l/25g body weight of Aurovist, Nanoprobe, USA)  
119  $\mu$ CT scan (Triumph II, TriFoil Imaging, USA) provided the geometrical information of the  
120 carotid bifurcation at an isotropic resolution of 50  $\mu$ m.

### 121 Post-processing of the measurements

122 The outer envelope of the Doppler signal and the diameter distension of the M-mode meas-  
123 urements were tracked to obtain respectively the blood velocity and the diameter curve. For  
124 the three locations (common, external (ECA) and internal (ICA) carotid artery) both signals  
125 were aligned in time. In order to calculate the instantaneous average velocity over the cross-  
126 section, the maximum velocity was divided by 2 assuming a parabolic flow profile at the  
127 measurement location (which was confirmed a posteriori by the simulations). Due to meas-  
128 urement errors and assumptions in the data processing, the time-averaged mass flow balance  
129 in the model (cycle-averaged difference between in- and outflow) was not completely fulfilled.

130 As we estimated the CCA measurement to be the most accurate, a scaling factor  $c=1.1275$   
131 was applied to the ICA and ECA mass flows such that the measured mass flow at the CCA  
132 was equal to the sum of the measured mass flows at ICA and ECA, multiplied with  $c$ . Mimics  
133 (Materialise, Leuven, Belgium) was used to semi-automatically segment the three branches of  
134 the carotid bifurcation from the  $\mu$ CT datasets. Flow extensions with a length equal to the di-  
135 ameter of the arteries were added in VMTK ([www.vmtk.org](http://www.vmtk.org)). Finally, the segmented geome-  
136 try was shrunk to the diastolic radius as measured with M-mode, resulting in a geometrical  
137 shrinking factor of 0.84. From the resulting STL-surface, a volume mesh of both the fluid  
138 (lumen) and the solid (arterial wall) domain was constructed using the in-house developed  
139 XTM meshing method (Bols et al., 2016). A mesh convergence study resulted in convergence  
140 for 300k fluid domain cells and 65k solid domain cells (De Wilde et al., 2015a). The wall  
141 thickness of the artery was assumed to be 10% of the local luminal diameter.

#### 142 **Steady state CFD simulations**

143 For the steady state CFD simulations 2 sets of BCs were applied: (i) the average velocity/flow  
144 of the three branches and (ii) the peak velocity of the three branches. For both cases, the mass-  
145 flow of the three branches was calculated as described above. At the common carotid the cal-  
146 culated mass flow was applied. At the external and internal carotids the mass flow split was  
147 calculated and they were imposed as outflow boundary conditions with a constant outflow  
148 fraction. For the steady state simulations convergence was reached when the continuity and  
149 momentum residuals dropped below  $1e^{-10}$ .

#### 150 **Transient CFD simulations**

151 For the transient CFD simulations a mass flow inlet was applied at both the common and ex-  
152 ternal carotid (at the external carotid the mass flow inlet behaved as an outflow). The internal

153 carotid was modelled as a traction free outlet. For the transient CFD simulations a conver-  
154 gence criterion of  $1e^{-10}$  was applied for both the continuity and the momentum equations.

### 155 **Transient FSI simulations**

156 The methodology for the FSI simulations is described in depth in our previously published  
157 work (De Wilde et al., 2015a). In short, the same mass flow inlets as for transient CFD simu-  
158 lations were applied. At the ICA a 3-element windkessel model was implemented, with pa-  
159 rameters fitted to the mouse-specific ultrasound measurements. The external tissue support  
160 was modelled using linear springs. The backward incremental method was executed sequen-  
161 tially to take axial and diastolic pre-stresses into account (De Wilde et al., 2015a). The abso-  
162 lute coupling iterations convergence criteria for the residuals were set to  $10^{-8}$  m for the  
163 displacement and 0.5 Pa for the load at the interface.

### 164 **Influence of locomotor activity**

165 During a period of high locomotor activity (=active state), the hemodynamic parameters in-  
166 crease even further than in the conscious resting case. The reported increase is: heart rate +  
167 23.77%, blood flow +93% and blood pressure +31.3% (Janssen et al., 2004). This extra in-  
168 crease rescaled case *Rest1* to case *Active1*, the lower bound estimate of the active condition.  
169 Similarly, *Active2* was defined as the adaptation of case *Rest2* with regard to the influence of  
170 locomotor activity (Table 1). In both active cases the heart rate, blood flow and blood pressure  
171 are increased substantially. Once again, the change in heart rate could only be applied to the  
172 transient simulations (transient CFD and FSI) and the blood pressure only for the FSI simula-  
173 tions. For the CFD simulations, the boundary conditions were simply scaled according to the  
174 rescaling factors of Table 1. For the FSI simulations the first step was to recalculate the 3-  
175 element windkessel model terminating the ICA based on the new, rescaled flow and blood

176 pressure. It is important to note that the rescaled pressure was not applied explicitly, but im-  
177 plicitly through this windkessel approach. For both the common and external carotids the re-  
178 scaled mass flows were applied at the boundaries. The pre-stress calculation was not repeated,  
179 because the previously found diastolic stress-state (at a pressure of 87 mmHg) was still valid.

## 180 **Post-processing of simulated data**

181 WSS was calculated for steady-state CFD simulations, and TAWSS for transient CFD and  
182 FSI simulations. WSS and TAWSS were normalized by their mean value to allow for a com-  
183 parison of patterns rather than values. For transient simulations the OSI was calculated as well  
184 (already dimensionless). The influence of anesthesia was visualized by the diastolic state of  
185 FSI simulations. In order to visualize the presence of recirculation structures, the line integral  
186 convolution (LIC) of the normalized WSS was plotted in the early deceleration phase. The  
187 LIC is the mathematical equivalent on a surface of putting drops of ink in a flow/vector field  
188 and was calculated for FSI simulations at different activity levels. Finally, the area experienc-  
189 ing reversed WSS was calculated for each computational technique and for each boundary  
190 condition. Reversed WSS was quantified using the scalar product of the WSS with the normal  
191 of the CCA inlet boundary. If the scalar product was negative (component in the direction of  
192 the CCA), the zone was indicated to have a reversed WSS. For the transient simulations the  
193 calculation of reversed WSS was done in the early deceleration phase. This definition of re-  
194 versed WSS was not applicable when the WSS vector was perpendicular to the inlet normal,  
195 which was the case at the stagnation point of the bifurcation.

## 196 **Statistics**

197 In order to compare different simulation techniques and activity levels we adopted a statistical  
198 technique analysis that has been described in detail in our previous work (De Wilde et al.,  
199 2015b). Briefly, we applied the surrogate sample data analysis (Peiffer et al., 2013a) to study

200 correlations between the calculated hemodynamic wall parameters obtained from different  
201 simulations. The similarity of the distribution of two different variables is described by  
202 Spearman's rank correlation coefficient  $\rho$  (Rowland et al., 2015) and Pearson's correlation  
203 coefficient. The correlation is determined for the original values of the variables and for  $n_{\text{sur}}$   
204 scrambled distributions of one of the variables (=surrogates). The statistical test ascertains  
205 that the original correlation coefficient is higher than the 95% percentile of the random distri-  
206 butions (i.e. the right-hand side of the histogram – right tailed test). To take into account the  
207 spatial autocorrelation of the variables, extra sampling steps ( $n_{\text{smp}}$ ) were performed. The sam-  
208 ple size  $s_{\text{smp}}$  was calculated for every geometry using:  $s_{\text{smp}} = s_{\text{map}}/l_{\text{dec}}^2$ . In this formula  $s_{\text{map}}$  was  
209 the original number of grid points and  $l_{\text{dec}}$  the spatial de-correlation length. A conservative  
210 value of 20 was used for  $l_{\text{dec}}$  in all the surrogate sample data analysis calculations to assure  
211 that all sample points were independent observations and that a proper statistical testing was  
212 conducted. The final histogram was averaged over  $n_{\text{smp}}$  histograms (one for each sampling  
213 step), with confidence levels as shown in figure 2. An error tolerance of  $10^{-4}$  was used for the  
214 p-value resulting in on average:  $n_{\text{sur}} = 2600$  and  $n_{\text{smp}}=3300$ . This technique was used to calcu-  
215 late the correlation between different simulations. The found correlation was deemed signifi-  
216 cant at  $p<0.05$  if the lower confidence bound (LCB)  $>0$  for a right tailed test (expected  $\rho>0$  –  
217 other HWP). All statistics were done on a bifurcation level, combining the three branches into  
218 a single statistical test.

219

220

221

## 222 RESULTS

### 223 Comparison of computational techniques

224 The flow-split going to the internal and external carotid branches differed between the mean  
225 flow steady state CFD simulation (0.13/0.87) and the peak flow steady state CFD simulation  
226 (0.24/0.76). The transient CFD flow in the internal carotid (Figure 1a) was much more peaked  
227 compared to the corresponding transient FSI flow (Figure 1b): the FSI simulations approxi-  
228 mated the measurements markedly better than the CFD simulations. The (TA)WSS calculated  
229 by all three CFD techniques (mean steady state CFD, transient CFD and transient CFD)  
230 reached a high and significant correlation ( $r=0.99$ ,  $0.96$  and  $0.99$ ;  $\rho=0.99$ ,  $0.94$  and  $0.99$  re-  
231 spectively) with the TAWSS calculated by FSI (Figure 2a). Nevertheless WSS was substan-  
232 tially higher when the peak velocity was used as a boundary condition for steady state CFD  
233 simulations (Figure 3a). After normalization by their mean values ( $3.25$ ,  $7.54$ ,  $3.42$  and  $2.85$   
234 Pa for mean CFD, peak CFD, transient CFD and FSI respectively) the WSS patterns showed  
235 very similar distributions for all techniques (Figure 3b). Nevertheless, the transient CFD  
236 simulations showed an important difference in oscillatory behavior when compared to FSI.  
237 The correlation was no longer significant and dropped to  $r=0.03$  and  $\rho=0.30$  (Figure 2b). In  
238 the transient CFD simulation the OSI was negligible with an average value of  $1.3 \cdot 10^{-3}$  and a  
239 maximum value of  $0.166$  (at the flow stagnation point, Figure 3c). For the FSI simulation on  
240 the other hand, there was a clear increase in OSI at the outer side of the ECA sinus, increasing  
241 the average value to  $2.46 \cdot 10^{-2}$  with a maximum value of  $0.272$ . FSI simulations also resulted  
242 in a larger area experiencing reversed WSS (Figure 4d).

### 243 **Influence of anesthesia and locomotor activity**

244 For all 4 increased activity states (*Rest1*, *Rest2*, *Active1* and *Active2*) the FSI-based TAWSS  
245 was highly and significantly correlated to the TAWSS calculated under anesthesia (Figure 2c,  
246  $r=0.99$ ,  $0.99$ ,  $0.99$ ,  $0.97$  and  $\rho=0.99$ ,  $0.99$ ,  $0.98$ ,  $0.98$  respectively). The TAWSS increased  
247 with a rising activity level (Figure 4a), mainly due to the increase in flow. The normalized  
248 TAWSS distribution was, however, very similar for all simulated cases (Figure 4b). For *An-*  
249 *aesthesia*, *Rest1*, *Rest2*, *Active1* and *Active2* the average TAWSS was 2.85, 3.08, 5.12, 4.77  
250 and 8.11 Pa respectively. The most apparent difference between different states was the rela-  
251 tively high normalized TAWSS at the inner bend of the bifurcation. Unlike the normalized  
252 TAWSS, the OSI was markedly affected by an increase in activity level (Figure 2d). The cor-  
253 relation with the OSI calculated under anesthesia remained significant but the correlation co-  
254 efficients decreased slightly for increasing activity levels ( $r=0.99$ ,  $0.97$ ,  $0.97$ ,  $0.86$  and  
255  $\rho=0.99$ ,  $0.98$ ,  $0.94$ ,  $0.91$  respectively). Most of this discrepancy could be related to an increase  
256 in OSI near the ostium of the ECA with rising activity levels (Figure 4c). No difference in  
257 OSI was discernible in the CCA or the ICA. A detailed quantification of the surface area ex-  
258 perencing WSS reversal confirmed that reversed WSS occurs at the outer bend of the ECA  
259 sinus for cases *Rest2*, *Active1* and *Active2* (Figure 4d). Especially for the case *Active 2* a sub-  
260 stantial zone with reversed WSS was observed. The LIC confirmed that flow recirculations  
261 were most apparent in the cases with highest locomotor activity (Figure 5a). For the FSI simu-  
262 lation of *Active2*, the most extreme case with high locomotor activity and the highest estimat-  
263 ed influence of anesthesia, the LIC was calculated at different time points in the mid plane of  
264 the ECA (Figure 5b). This figure shows that the flow recirculation was maximal during the  
265 early deceleration phase ( $t_2$ ), which was also the time point used for Figure 5a. These data are  
266 consistent with the elevated OSI (Figure 4c) and the locally increased WSS reversal (Figure

267 4d) that were both observed at the outer bend of the ECA for case *Active2*. Due to the pres-  
268 ence of the flow recirculation the magnitude of the velocity vectors remained low at this loca-  
269 tion (Figure 5b).

270

271

## 272 **DISCUSSION**

### 273 **Comparison of computational techniques**

274 In this paper, several modelling strategies to simulate the murine carotid hemodynamics were  
275 compared. The focus of the comparison was on the WSS and other hemodynamic wall param-  
276 eters. We found that the (normalized) WSS pattern was virtually the same for all the model-  
277 ling strategies (Figure 2a, 3b), despite the clearly different boundary conditions that were  
278 applied (Figure 1a, 1b). Even the inclusion of the pulsatile nature of the blood flow (transient  
279 CFD and FSI) and/or the inclusion of the buffering capacity of the carotid artery (FSI) had  
280 little effect. Nonetheless, although the normalized WSS patterns were similar, the absolute  
281 values differed between CFD and FSI simulations because the total flow through the specific  
282 geometry was not constant and because the pulsatile pressure component led to a cyclic infla-  
283 tion of the arteries in the FSI simulations, which influenced the cross-sectional area (Figure  
284 1c). A larger impact was found for parameters reflecting the oscillatory behavior of the WSS  
285 (Figure 2, 3c, 4d). The clearly higher oscillatory character of the ECA flow in the FSI simula-  
286 tions is directly related to the buffering capacity of the latter. Thanks to this buffering capacity  
287 FSI allowed for a net instantaneous flow imbalance and thus a net in/outflow difference, while  
288 the CFD simulations had to comply with an instantaneous mass balance at all times. Our data  
289 therefore suggest that in the carotid artery of mice FSI offers less restrictive simulation condi-  
290 tions, and may capture small oscillations that were not present in CFD.

### 291 **Influence of anesthesia and locomotor activity**

292 Decreased flow measurements due to anesthesia can have a big influence on the simulation  
293 results (Mohri et al., 2014). Confirming literature data, we also observed the presence of a  
294 flow recirculation at the outer sinus of the ECA when conscious conditions and locomotor

295 activity were taken into account (Figure 4). As a result of the changed boundary conditions  
296 (Table 1) the peak systolic Reynolds number in the common carotid changed from 13.38 (*An-*  
297 *esthesia*) to 39.57 (*Active2*). For the same cases the Womersley number increased from 0.44  
298 to 0.70, indicating the possibility of changes in flow behavior. These data indicate that one  
299 should interpret the results of simulations that assume anesthesia with caution. Our normal-  
300 ized TAWSS patterns were rather robust with regard to different activity levels. However, ab-  
301 solute values of the WSS/TAWSS were heavily influenced by the precise boundary  
302 conditions and the level of physical activity that were assumed in the simulations. Therefore it  
303 makes sense for future work to report normalized WSS metrics rather than absolute values.  
304 Finally, the OSI (Figure 4c) and the reversed WSS (Figure 4d) showed a markedly different  
305 distribution depending on the physiological state. The location of the region with elevated  
306 OSI and reversed WSS was similar for all simulations, but the shape and size of this location  
307 was not, with the largest increase in the presence of extensive locomotor activity. These data  
308 suggest that flow recirculation may develop in highly active mice while it is not present under  
309 anesthesia or conscious resting conditions.

### 310 **Implications for the interpretation of previously published work**

311 We believe that the largest impact of our findings lies in the (re-)interpretation of CFD simu-  
312 lations in mice. Many authors have tried to link local disease patterns to focal zones of locally  
313 disturbed shear stress using rigid-walled models of either the aortic arch (Assemat et al., 2014;  
314 Feintuch et al., 2007; Suo et al., 2007; Yap et al., 2014), the abdominal aorta (Ford et al.,  
315 2011; Greve et al., 2006; Trachet et al., 2011b; Willett et al., 2010) or both the thoracic and  
316 abdominal aorta (Hoi et al., 2011; Huo et al., 2008; Van Doormaal et al., 2014). The findings  
317 of publications that focused on regular WSS patterns (Feintuch et al., 2007; Greve et al., 2006;  
318 Huo et al., 2008; Suo et al., 2007; Willett et al., 2010; Yap et al., 2014) are not affected much

319 by our results. However, the findings of publications that focused on the oscillatory patterns  
320 of WSS (Assemat et al., 2014; Ford et al., 2011; Hoi et al., 2011; Trachet et al., 2011b; Van  
321 Doormaal et al., 2014) might need to be re-interpreted to include distending aortic walls. On  
322 the other hand, the fact that most of these CFD simulations were based on boundary condi-  
323 tions measured under anesthesia does not seem to compromise their findings too much. The  
324 latter remark also holds for publications that estimated shear stress patterns in the carotid ar-  
325 tery from Doppler measurements under anesthesia. In their seminal paper from 2006, Cheng  
326 et al used this technique to demonstrate that atherosclerotic lesions developed invariably in  
327 the regions of the carotid artery where vortices and low shear stress had been induced by a  
328 perivascular shear stress modifier, whereas the regions in which the device induced an in-  
329 creased shear stress were protected from plaque development (Cheng et al., 2006). While the  
330 perivascular cast may have induced more vortices in conscious mice than what was previously  
331 assumed under anesthesia, our results do not change the basic findings of this paper.

### 332 **Limitations and future work**

333 Our simulations are inherently limited by the spatial and temporal resolution of the preclinical  
334 imaging that provided the geometry and boundary conditions. But while future improvements  
335 in imaging might improve the level of details (e.g. branches etc.) that can be included  
336 (Trachet et al., 2015), we are confident that the main message of this paper will not be affect-  
337 ed. It should be noted that our estimated increase in blood flow through the common carotid  
338 artery due to locomotor activity was probably on the high side and thus represents a worst  
339 case scenario. At high locomotor activity, the cardiac output increases because of an increase  
340 in blood flow towards the skeletal muscles, while the main perfusion area of the carotid arter-  
341 ies is the brain (Marieb and Hoehn, 2008). In humans, the cerebral blood flow only peaks at  
342 the onset of physical activity and decreases again during constant physical activity, while the

343 cardiac output is still at an elevated level (Querido and Sheel, 2012). The initial increase due  
344 to physical exercise (steady-state cycling) has been quantified for healthy humans as  $27.9 \pm$   
345  $28.6\%$  and the increase during prolonged exercise has been reported  $2.6 \pm 13.5\%$  (Hiura et al.,  
346 2014). Both of these changes are smaller than the increase in cardiac output which is 92% and  
347 118% respectively. It is thus likely that when mice are active the change in blood flow  
348 through the carotid arteries is smaller than the change in cardiac output. If the blood flow to  
349 the brain would be assumed to remain constant during (prolonged) locomotor activity, cases  
350 *Active1* and *Active2* should be ignored and *Rest2* would be the most extreme case. *Rest2* is  
351 exactly at the level where the recirculation started developing in this geometry. In conclusion,  
352 our data suggest that it is challenging to estimate the effect of anesthesia and locomotor activi-  
353 ty correctly, and more research is needed to improve the accuracy of the corresponding  
354 boundary conditions.

355

356 **CONCLUSIONS**

357 Simple steady state CFD simulations can be used to assess the WSS pattern in the murine ca-  
358 rotid bifurcation. However, the average WSS value differs between simulation strategies. To  
359 be able to assess the oscillatory behavior of the WSS (OSI, reversed WSS), the buffering ef-  
360 fect of the elastic carotid artery should be taken into account and FSI simulations are indis-  
361 pensable. The effect of anesthesia and locomotor activity on the WSS patterns is minimal, but  
362 flow recirculation and WSS reversal that are not present under anesthesia do occur when the  
363 mice are in a high locomotor activity state.

364

365 **ACKNOWLEDGMENTS**

366 David De Wilde is supported by a research grant of the Flemish government agency for Inno-  
367 vation by Science and Technology (IWT, Grant number 111618). Bram Trachet is supported  
368 by a research grant of the Research Fund - Flanders (FWO, Grant number 12A5816N). We  
369 thank Christian Vanhove, Benedicte Descamps and Scharon Bruneel for their assistance with  
370 the in vivo experiments, Francisco Londono and Abigail Swillens for their assistance in pro-  
371 cessing the ultrasound measurements, Bert Vandeghinste for his assistance in processing the  
372 micro-CT scans and Mathias Peirlinck, Liesbeth Taelman, Joris Bols, Francesco Iannaccone,  
373 Nic Debusschere, Joris Degroote and Jan Vierendeels for their assistance with the numerical  
374 simulations.

375 **CONFLICT OF INTEREST STATEMENT**

376 No conflict of interest

377

378

379

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479

480

481 **TABLES**

482

483

<b>Scale factor/ Condition</b>	<b>Q</b>	<b>P<sub>sys</sub></b> (mmHg)	<b>P<sub>dia</sub></b> (mmHg)	<b>T</b> (s)	<b>HR</b> (beats/min)
<b>Anesthesia</b>	1.00	119	87	0.1719	349
<b>Rest1</b>	1.06	119	87	0.1530	392
<b>Rest2</b>	1.72	119	87	0.1066	563
<b>Active1</b>	2.04	156	114	0.1236	485
<b>Active2</b>	3.33	156	114	0.0861	697

484

485 Table 1. Adaptation of the boundary conditions to account for anesthesia and locomotor activ-

486 ity.

487

488 **FIGURE LEGENDS**

489

490 **Figure 1.** Imposed flow profiles for transient CFD (a) and FSI (b) simulations, and resulting  
491 cross-sectional areas (c) at the in-and outlets of the model.

492

493 **Figure 2.** (a) Correlation between the TAWSS as calculated by FSI (assuming anesthesia) and  
494 the (TA)WSS calculated by mean, peak and transient CFD. (b) Correlation between the OSI  
495 as calculated by FSI (assuming anesthesia) and the OSI calculated by transient CFD. (c) Cor-  
496 relation between the TAWSS calculated assuming anesthesia and the TAWSS calculated as-  
497 suming different locomotor activity levels. Remark that the Y-axis is not on the same scale, as  
498 the absolute value of TAWSS increased with increasing activity (and flow). (d) Correlation  
499 between the OSI calculated assuming anesthesia and the OSI calculated assuming different  
500 locomotor activity levels. \*:  $p < 0.05$ , \*\*:  $p < 0.001$ ; ns: not significant.

501

502 **Figure 3.** Qualitative comparison of the influence of computational techniques (assuming an-  
503 esthesia) for TAWSS (a), normalized TAWSS (b) and OSI (c).

504

505 **Figure 4.** Qualitative comparison of the influence of anesthesia and locomotor activity (using  
506 FSI simulations) on TAWSS (a), normalized TAWSS (b), OSI (c). In panel d the reversed  
507 WSS area is quantified for all techniques and all physiological conditions. The close-up  
508 shows the transient CFD simulation with both reversed WSS zones: stagnation point (blue)  
509 and the outer sinus of the ECA (red).

510

511 **Figure 5.** The line integral convolution (LIC) shows recirculation zones in the WSS at the  
512 anterior side of the ECA (a) and in the mid-plane of the ECA (b). The LIC is plotted at the  
513 early deceleration phase for different activity states (a), and at different time points for the  
514 *Active2* case (b).

515

516

517