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# CFD-modelling of activated carbon fibers for indoor air purification

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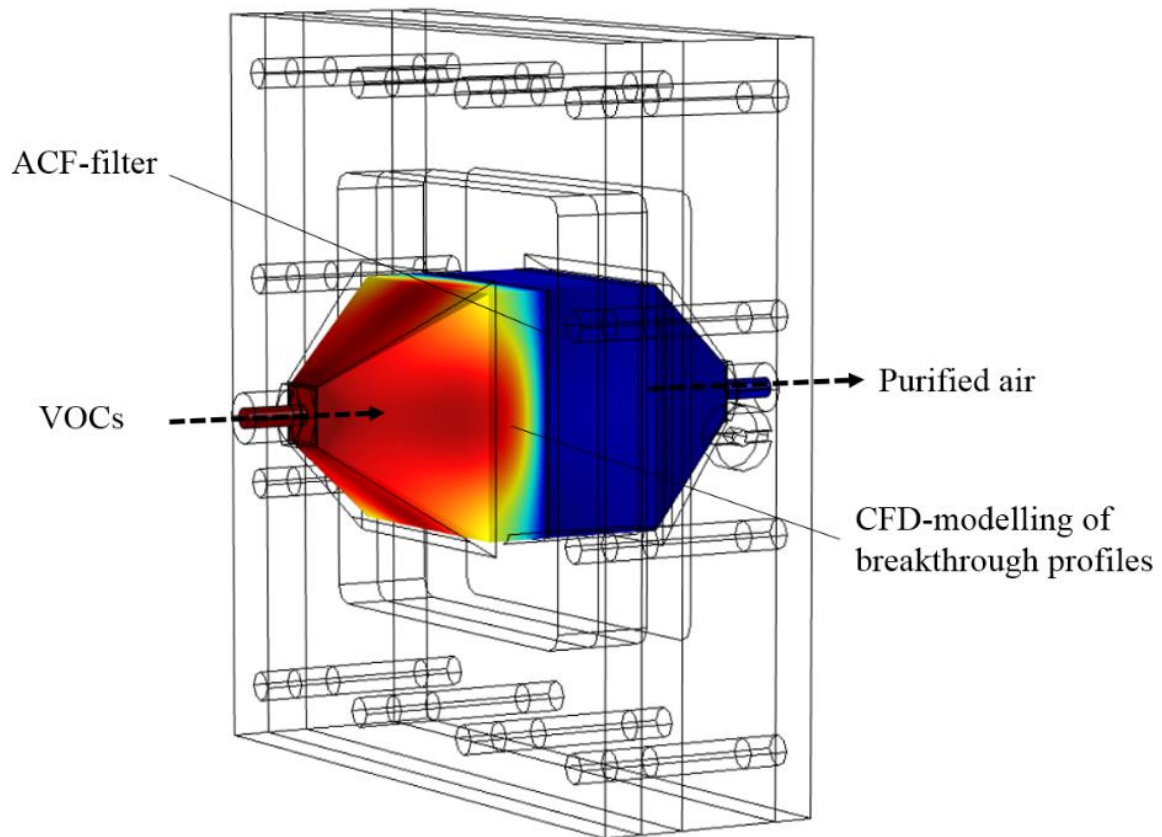
## Keywords

Computational Fluid Dynamics; Activated Carbon Fiber; Pressure drop; Adsorption isotherm;  
Breakthrough profiles

## Highlights

- Pressure drop measurements for various activated carbon fiber filter thicknesses.
- CFD simulations and fitting of ACF pressure drop with Darcy-Forchheimer law.
- Investigation of three adsorption isotherms to model dynamic adsorption behavior.
- Successful Multiphysics simulation of adsorption breakthrough curves of acetaldehyde

# 1 Graphical Abstract



2

## 3 1. Abstract

4 Activated carbon fibers for indoor air purification were investigated by means of pressure drop  
5 and adsorption capacity. The Darcy-Forchheimer law combined with Computational Fluid  
6 Dynamics (CFD) modelling was deployed to simulate the pressure drop over an activated  
7 carbon fiber (ACF) filter with varying filter thickness. The CFD model was later combined with  
8 adsorption modelling to simulate breakthrough profiles of acetaldehyde adsorption on the ACF-  
9 filter. The adsorption model incorporates mass transfer resistance and adsorption equilibrium.  
10 It assumes local equilibrium between gas phase and solid phase. The latter was investigated for  
11 three different adsorption isotherms: linear, Langmuir and Freundlich adsorption. Successful  
12 agreement between model simulations and experimental data was obtained, using the

1 Freundlich adsorption model. The numerical model could provide valuable insights and allows  
2 to continuously improve the design of filtration devices.

3

## 4 **2. Introduction**

5 Indoor air quality has become an issue of major interest due to stringent insulation and  
6 airtightness measures in buildings [1]. Accumulation of volatile organic compounds (VOCs)  
7 can be counteracted by ventilation with outdoor air, but this demands a high energy  
8 consumption due to increased ventilation rates and conditioning of the outdoor air. Sidheswaran  
9 et al. state that a HVAC system, operating at 15-20% outdoor air and 80-85% indoor air is  
10 viable, provided that 15 to 20% of the VOCs are removed in the supply air flow [2]. Among the  
11 different strategies, activated carbon filters have been suggested as a promising approach for  
12 effective VOC removal [3–5]. Activated carbon is a widely used solution for cleaning industrial  
13 waste streams with high VOC concentrations, but recent studies have shown that activated  
14 carbon is also feasible for purification in building applications, with VOC concentrations at ppb  
15 levels [6–8]. For the same operational conditions, activated carbon fiber cloths or felts offer  
16 many advantages over conventional granular activated carbon, such as higher adsorption  
17 capacities, faster diffusion rates and kinetics, lower pressure drop, more compactness, great  
18 flexibility and easier regeneration [2,4,7,9–13]. Regeneration of activated carbon fibers can be  
19 accomplished by transmitting an electrical current, inducing the joule-effect and hence, heating  
20 the filter and desorbing the pollutants [14]. Although this technique seems very promising, there  
21 is still some investigation required to develop a commercial device for implementation in  
22 Heating, Ventilation and Air Conditioning (HVAC) systems. Integration of an ACF-filter into a  
23 HVAC system requires special features to deal with the typical high flow rates. The filter must  
24 allow the polluted air to pass through with a minimal pressure drop, but still feature sufficient  
25 adsorption capacity to effectively remove the VOCs. In this work, we investigated the pressure

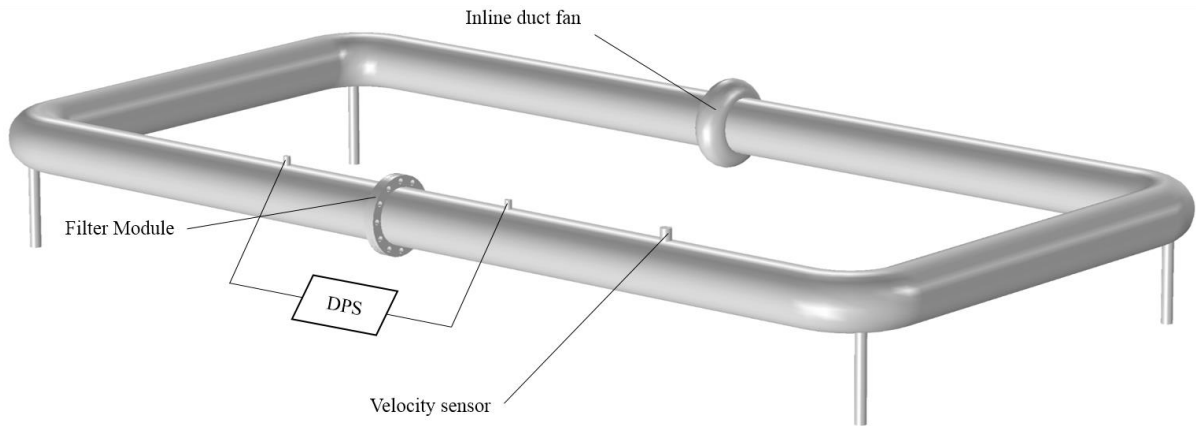
1 drop over an activated carbon fiber cloth with varying filter thickness for typical HVAC flow  
2 rates. In addition, we developed a numerical model to validate the Darcy-Forchheimer  
3 parameters, which were calculated based on the experimental pressure curves, and to simulate  
4 the transient adsorption behavior and breakthrough curves of air contaminated with  
5 acetaldehyde, as a model compound for VOCs in indoor air. Adsorption of acetaldehyde was  
6 fitted with experimental data by three common adsorption models: Linear, Langmuir-type and  
7 Freundlich-type adsorption models. Several studies have shown that these adsorption models  
8 are valid for adsorption of volatile organic compounds on ACF filters [4,5,7,15] and some  
9 research has been conducted to model dynamic adsorption on activated carbon fibers, based on  
10 these adsorption models [7,16]. In this work, a CFD analysis is presented, combining complex  
11 fluid dynamics with adsorption kinetics to gain more valuable insights in the adsorption process.  
12 Such a model is very useful to evaluate the performance of an ACF system in unsteady operating  
13 conditions (typical for realistic applications) and can be used as a tool to develop commercial  
14 devices that can be integrated in HVAC systems [4,17,18].

### 15 **3. Methodology**

#### 16 **3.1 Pressure curves**

17 For the development of an efficient ACF filtering device, a trade-off has to be made between  
18 the number of filter layers to increase the adsorption capacity and a maximum pressure drop for  
19 energy consumption reasons, which is evidently increased with the number of layers. It is  
20 therefore important that the filter thickness (and thus adsorption capacity) is adjusted to the  
21 (maximum) concentration of VOC in the supplying air stream. In the first part of this work, the  
22 pressure drop across an FM10 (Calgon Corp) ACF filter is studied in a closed circuit air duct  
23 system for varying thicknesses of the filter (0.5-2 mm) at typical HVAC flow rates up to 2 m/s  
24 ( $T=20^{\circ}\text{C}$ ,  $P=101.3\text{ kPa}$ ). The FM10 woven cloth ACF filter has been studied in the past in terms

1 of surface density, specific surface area, apparent density, particle radius, porosity, etc. [2,7,8].  
2 The setup consists of circular pipes with an inner diameter of 103 mm and a total length of 6m



**Figure 1** Closed circuit vent duct with inline duct fan, velocity sensor and differential pressure sensor (DPS)

3 (Figure 1). Up to four layers of ACF filter were compressed between two polylactic rings and  
4 the module was placed in line with the pipe duct. This way, the air flows perpendicular to the  
5 surface area of the filter in every experiment. The air flow rate was controlled by an inline duct  
6 fan (Ruck RS-series) with a free flow rate of 350 m<sup>3</sup>/h and a static pressure of 250 Pa. The fan  
7 was mounted directly on the external rotor motor with controllable voltage (24V-230V) through  
8 which the air flow rate could be adjusted. The air flow rate was measured by a hotwire air  
9 velocity sensor (CTV 110, KIMO instruments) in the middle of the vent duct and the pressure  
10 drop was measured by a differential pressure sensor (DPS) consisting of a pressure module  
11 (Fluke 717 30G, Fluke Corp.) and a pressure calibrator (Fluke 700PD2, Fluke Corp.). The air  
12 flow profile was assumed to be uniform, as this was investigated earlier by Koch et al. for the  
13 same vent duct [19].

14 Consequently, the pressure-velocity curves were used to determine Darcy-Forchheimer  
15 parameters (eq. 1), i.e. the permeability  $\kappa$  (m<sup>2</sup>) of the medium, which represents the flow  
16 resistance of the filter and the Forchheimer term  $\beta$  (1/m), which accounts for the non-linear  
17 behavior of the pressure difference at high flow rates due to inertial effects:

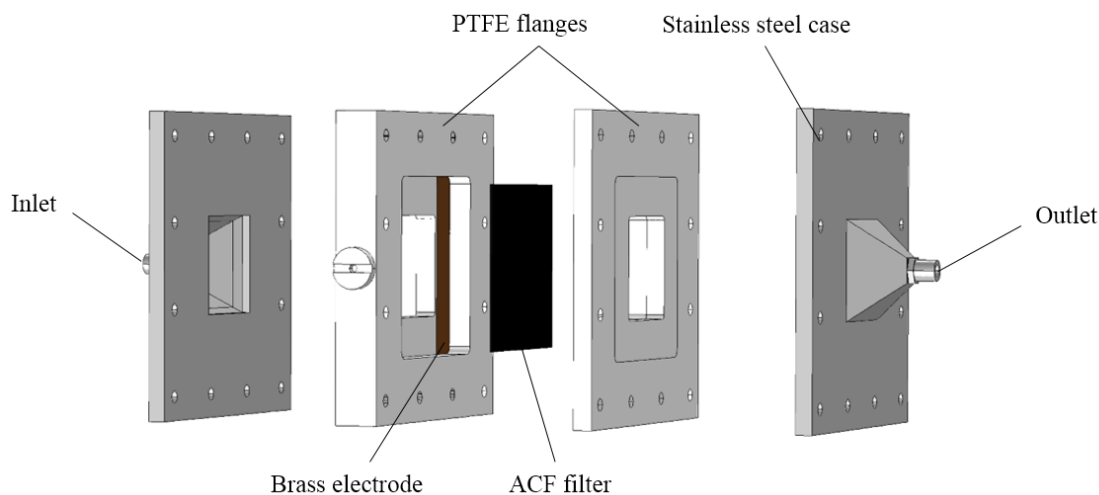
$$\frac{\Delta P}{\Delta x} = -\frac{\mu}{\kappa}v - \beta\rho v^2 \quad 1$$

1 With  $\Delta P$  the pressure drop (Pa),  $\Delta x$  de filter thickness (m),  $\mu$  the dynamic viscosity of the air  
2 flow (Pa·s),  $\rho$  the density of the air flow (kg/m<sup>3</sup>) and  $v$  the air flow velocity (m/s). Since the  
3 Darcy-Forchheimer law accounts for the filter thickness  $\Delta x$ , the parameters  $\kappa$  and  $\beta$  are  
4 expected to be similar for every experiment. These parameters were later validated by numerical  
5 simulation with Comsol Multiphysics.

6

### 7 **3.2 Adsorption experiments**

8 The adsorption capacity was investigated using acetaldehyde (Messer, 1% in N<sub>2</sub>) as a model  
9 compound for indoor air pollution. Acetaldehyde is more difficult to adsorb onto activated  
10 carbon than for example toluene due to its higher polarity, making it an ideal component to  
11 determine the minimum required adsorption capacity under given operating conditions. A filter  
12 module was developed based on the pilot unit described in Subrenat et al. (2001). The module  
13 is composed of 2 stainless steel cases, containing 2 PTFE flanges to ensure thermal and  
14 electrical insulation during the regeneration process (Figure 2). The cases were formed  
15 conically to gradually diffuse and constrict the air flow between the inlet or outlet (4 mm  
16 diameter) and the section where the carbon filter was fixed (50 × 50 mm flow-through surface).  
17 One of the PTFE flanges was equipped with two brass plate electrodes in order to provide the  
18 current for regeneration. One to four layers of carbon filter (5cm x 5cm) rested on the two  
19 electrodes, and were pressed between the two PTFE flanges. The two electrodes were externally  
20 connected to a Tenma 72-10480 DC power supply (0-30V output). The FM10 filter with the  
21 smallest thickness of 0.5mm (1 layer) was used in all adsorption experiments.



**Figure 2** Configuration of the ACF-filter module with regeneration electrodes

1

2 Acetaldehyde in air was supplied to the filter module ( $T=20^{\circ}\text{C}$ ,  $P=101.2\text{kPa}$ ) by mass flow

3 controllers at a flow rate of  $200\text{ cm}^3/\text{min}$  and measured at the outlet, using FTIR spectroscopy

4 (IR peak height at  $2728\text{ cm}^{-1}$ , which corresponds to the  $\nu(\text{C-H})$  stretch vibration [20]. The

5 experiment continued until the breakthrough curve was complete, i.e. the outlet concentration

6 matched the inlet concentration. The adsorbed amount of acetaldehyde at equilibrium was

7 derived by integrating the transient acetaldehyde concentrations over the time period and

8 subtracting the obtained value from the amount of acetaldehyde entering, a similar method was

9 performed in previous work [21]. After each adsorption cycle, the filter was regenerated, using

10 an electrical current and a clean air flow of  $2000\text{ cm}^3/\text{min}$ . A constant current of  $2.75\text{ A}$  ( $20\text{V}$ )

11 was imposed to achieve a steady desorption temperature of  $150^{\circ}\text{C}$  due to the Joule effect, as

12 measured with a T-type thermocouple at the surface of the filter. This temperature was persisted

13 until acetaldehyde could not be detected anymore by the on-line FTIR measurements.

14 Integration of the desorption curve verified that the amount of desorbed acetaldehyde equaled

15 the amount of adsorbed acetaldehyde in the prior adsorption experiment. The breakthrough

16 curves for 5 different inlet concentrations of acetaldehyde were studied:  $1.61$ ,  $4.75$ ,  $9.74$ ,  $14.50$

17 and  $19.19\text{ mmol}/\text{m}^3$  with two repetitions per concentration.



### 1 3.3 Numerical model

2 The Darcy-Forchheimer parameters, derived from the experimental pressure-velocity curves,  
3 were validated with Computational Fluid Dynamics. The commercial software package Comsol  
4 Multiphysics v5.3 was used for finite element modelling of the air flow through the filter  
5 medium. A straight part of the closed circuit air duct system (Figure 1) was drawn in Comsol  
6 as a simple cylindrical pipe of length 200mm, including a small section of 0.5-2mm that  
7 represented the filter medium in the pipe. The air stream was simulated as an incompressible  
8 flow with the 'Free and Porous Media Flow' physics interface in Comsol. The filter was defined  
9 as a porous medium with a porosity  $\epsilon_p$ , permeability  $\kappa$  and Forchheimer drag  $\beta$ . The latter two  
10 were determined based on the pressure-velocity curves. The inlet condition was set to a uniform,  
11 normal flow velocity with the respective velocity flow rate of each experiment. A steady-state  
12 solution was obtained with a direct stationary solver, by simultaneously solving the mass  
13 continuity equation, the governing equations for the free flow ( $k$ - $\epsilon$ -turbulent model) and the  
14 equations for flow in porous media:

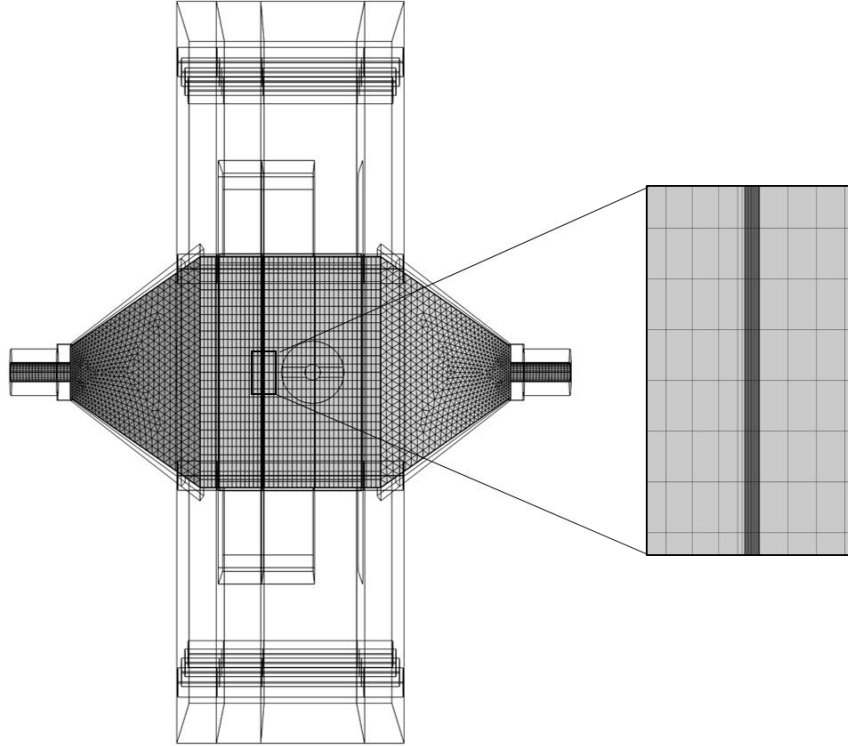
$$\rho \nabla(\mathbf{u}) = 0 \quad 2$$

$$\rho(\mathbf{u} \cdot \nabla)\mathbf{u} = \nabla \cdot (-p\mathbf{I} + (\mu + \mu_T)(\nabla\mathbf{u} + (\nabla\mathbf{u})^T)) \quad 3$$

$$\frac{\rho}{\epsilon_p} \left( (\mathbf{u} \cdot \nabla) \frac{\mathbf{u}}{\epsilon_p} \right) = \nabla \cdot \left( -p\mathbf{I} + \frac{\mu}{\epsilon_p} (\nabla\mathbf{u} + (\nabla\mathbf{u})^T) - \frac{2\mu}{3\epsilon_p} (\nabla \cdot \mathbf{u})\mathbf{I} \right) - (\mu\kappa^{-1} + \beta|\mathbf{u}|)\mathbf{u} \quad 4$$

15 With  $\rho$  the air density (1.2041 kg/m<sup>3</sup>),  $\mathbf{u}$  the velocity vector (m/s),  $\mathbf{I}$  the identity matrix,  $p$  the  
16 pressure (Pa),  $\mu$  the dynamic viscosity (Pa·s) and  $\mu_T$  the eddy viscosity (Pa·s), which is  
17 calculated by the  $k$ - $\epsilon$ -model for a turbulent flow. The last term in eq. 4 represents the Darcy-  
18 Forchheimer law. In order to accurately simulate the pressure drop in the filter, a very fine swept  
19 mesh was chosen for the porous medium with 30 elements in the flow direction. The total  
20 number of mesh elements for this geometry was 188,100 with a an average quality of 0.83. The

1 final aim of the numerical model was to simulate the breakthrough curves of the filter with  
 2 acetaldehyde as model compound. Therefore, the geometry of the filter module was constructed  
 3 and meshed in Comsol. The porous medium was meshed similarly as in the previous model and  
 4 the total mesh consisted of 562,400 elements (combined tetrahedral and hexahedral) with an  
 5 average quality of 0.72 (Figure 3).



**Figure 3** Mesh of the ACF-filter module

6

7 Transport of acetaldehyde in the porous medium was simulated by solving eq. 5:

$$\left( \epsilon_p + \rho_f \frac{\partial c_{ads}}{\partial c} \right) \frac{\partial c}{\partial t} + \left( c - \frac{\rho_f}{(1 - \epsilon_p)} c_{ads} \right) \frac{\partial \epsilon_p}{\partial t} + \nabla \cdot (-D_e \nabla c) + \mathbf{u} \cdot \nabla c = 0 \quad 5$$

8 With  $c$  the concentration of acetaldehyde ( $\text{mol/m}^3$ ),  $c_{ads}$  the amount of species adsorbed to the  
 9 solid phase ( $\text{mol/kg}$ ),  $\rho_f$  the filter density ( $\text{kg/m}^3$ ) and  $D_e$  the effective diffusivity within the  
 10 porous matrix ( $\text{m}^2/\text{s}$ ). The latter is determined by:

$$D_e = \frac{\epsilon_p}{\tau} D_K = \epsilon_p^2 D_K \quad 6$$

1 With  $\tau$  the tortuosity factor, approximated by  $1/\epsilon_p$ , and  $D_K$  the Knudson diffusion coefficient  
 2 (m<sup>2</sup>/s) [17,22].  $\frac{\partial c_{ads}}{\partial c}$  in eq. 5 is specified by the type of adsorption isotherm. Three adsorption  
 3 isotherm models were tested and validated against experimental data: Linear, Langmuir and  
 4 Freundlich adsorption, respectively expressed as:

$$c_{ads} = K_{lin} \cdot c ; \frac{\partial c_{ads}}{\partial c} = K_{lin}. \quad 7$$

$$c_{ads} = c_{ads,max} \frac{K_{Langm} \cdot c}{1 + K_{Langm} \cdot c} ; \frac{\partial c_{ads}}{\partial c} = \frac{K_{Langm} \cdot c_{ads,max}}{(1 + K_{Langm} \cdot c)^2} \quad 8$$

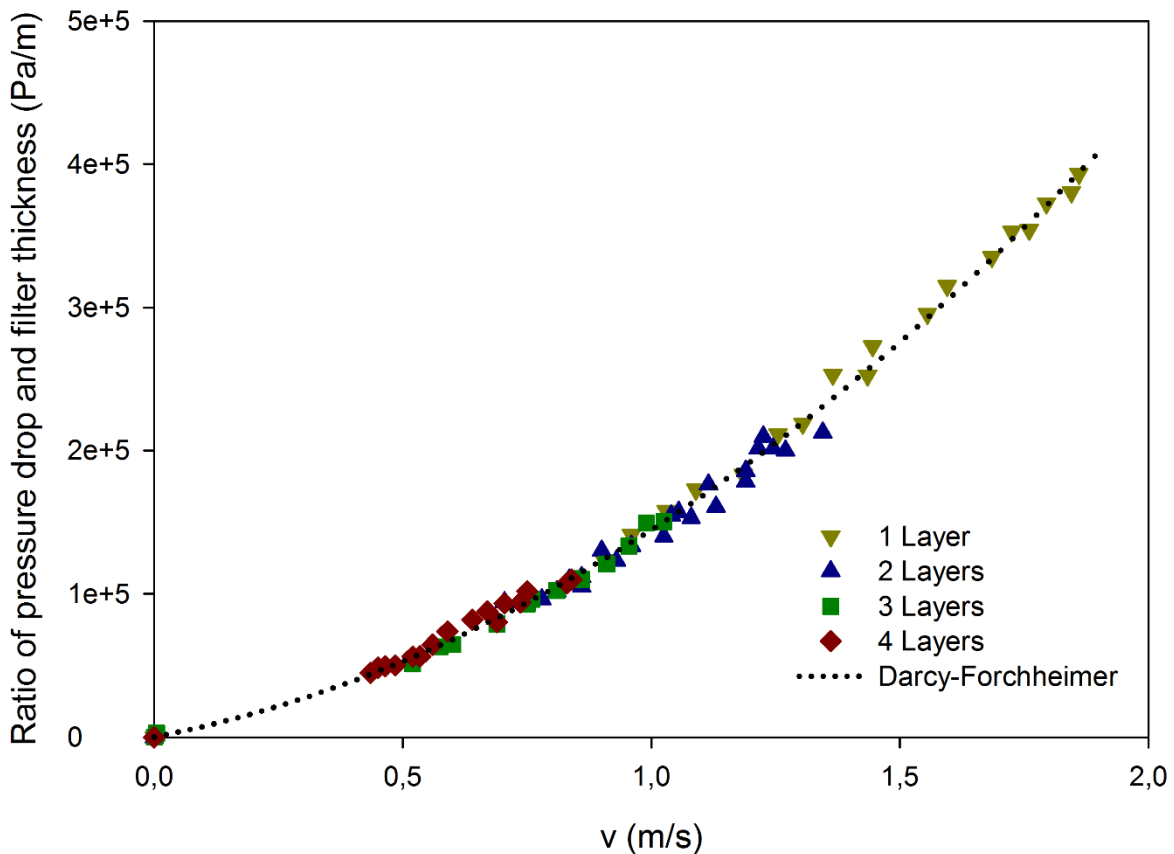
$$c_{ads} = K_{Freund} \cdot c^{1/n} ; \frac{\partial c_{ads}}{\partial c} = \frac{1}{n} \frac{c_{ads}}{c} \quad 9$$

5 Several studies have shown that Langmuir and Freundlich can effectively describe the  
 6 adsorption on activated carbon material [2,6,7,9,12,15]. However, in the case of Langmuir  
 7 adsorption, the expression is transformed to a linear adsorption when the concentration  $c$   
 8 (mol/m<sup>3</sup>) is very low, i.e.  $K_{Langm} \cdot c \ll 1$  [15]. The Freundlich isotherm can also be  
 9 approximated by a linear adsorption model when the exponent  $1/n$  approaches a value of 1.  
 10 The exponent  $1/n$  usually ranges from 0 to 1 for most species found in indoor air [4,18].  
 11 Therefore, we tested if these assumptions are valid, so that a simple linear model can be  
 12 employed to describe adsorption of acetaldehyde on the activated carbon filter. Firstly, a  
 13 parameter estimation was conducted for each adsorption model, based on the adsorbed amount  
 14 of acetaldehyde at equilibrium, as derived by integration of the breakthrough curves, and the  
 15 inlet concentration of acetaldehyde. Secondly, the estimated parameters were deployed in the  
 16 numerical model to simulate the breakthrough curves of different acetaldehyde concentrations.

## 1 4. Results and discussion

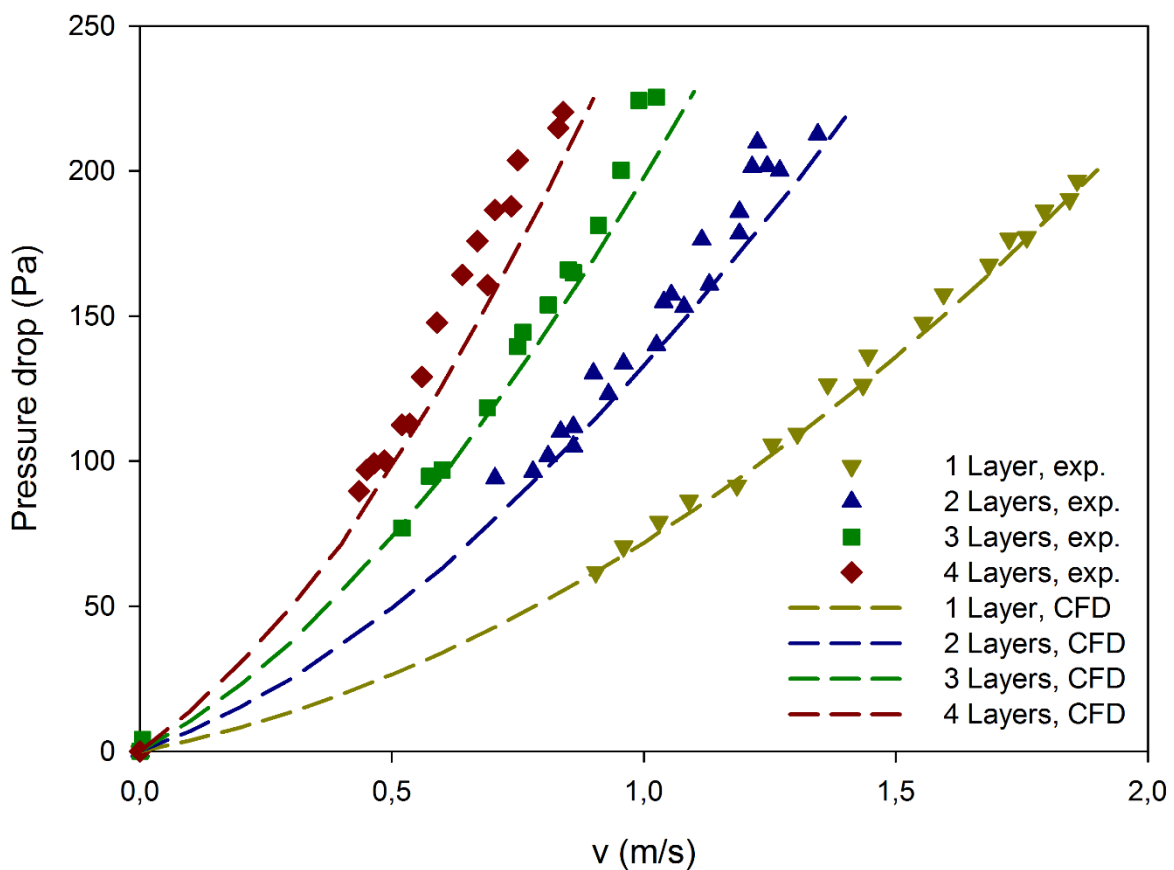
### 2 4.1 Pressure curves

3 Pressure curves were composed to determine the optimal filter thickness for typical HVAC flow  
4 rates. The pressure-velocity behavior at these flow rates can be represented best by the Darcy-  
5 Forchheimer law. In order to determine the Darcy-Forchheimer parameters, the data of all  
6 experiments with varying thicknesses were plotted, with the y-axis representing the ratio of  
7 pressure drop over filter thickness and the x-axis the air flow velocity (Figure 4). As seen from  
8 Figure 4, the data can be perfectly described by a polynomial expression of the second order,  
9 with a coefficient of determination of 0.995. From the quadratic regression, the permeability  $\kappa$   
10 and the Forchheimer term  $\beta$  could be resolved, i.e.  $2.71 \times 10^{-10} \text{ m}^2$  and  $64798 \text{ m}^{-1}$  respectively.



**Figure 4** Ratio of pressure drop and filter thickness (Pa/m) vs air flow velocities for various filter thicknesses. (0.5-2 mm) and the fitted Darcy-Forchheimer curve

1 Given the fact that the coefficient of determination approximates a value of 1, it can be assumed  
 2 that the Darcy-Forchheimer coefficients can be deployed to extrapolate the data to lower or  
 3 higher air flow velocities. Furthermore, the pressure drop over the ACF filter can perfectly be  
 4 simulated for a range of filter thicknesses while retaining the same set of parameters. In order  
 5 to validate these parameters, a CFD-model was developed in which the air flow was modeled  
 6 through a porous medium, with the latter representing the filter. The simulated pressure curves  
 7 are shown in Figure 5.



**Figure 5** CFD simulations of pressure drop versus air flow for individual filter thicknesses (0.5-2mm) and validation against experimental data

8

9

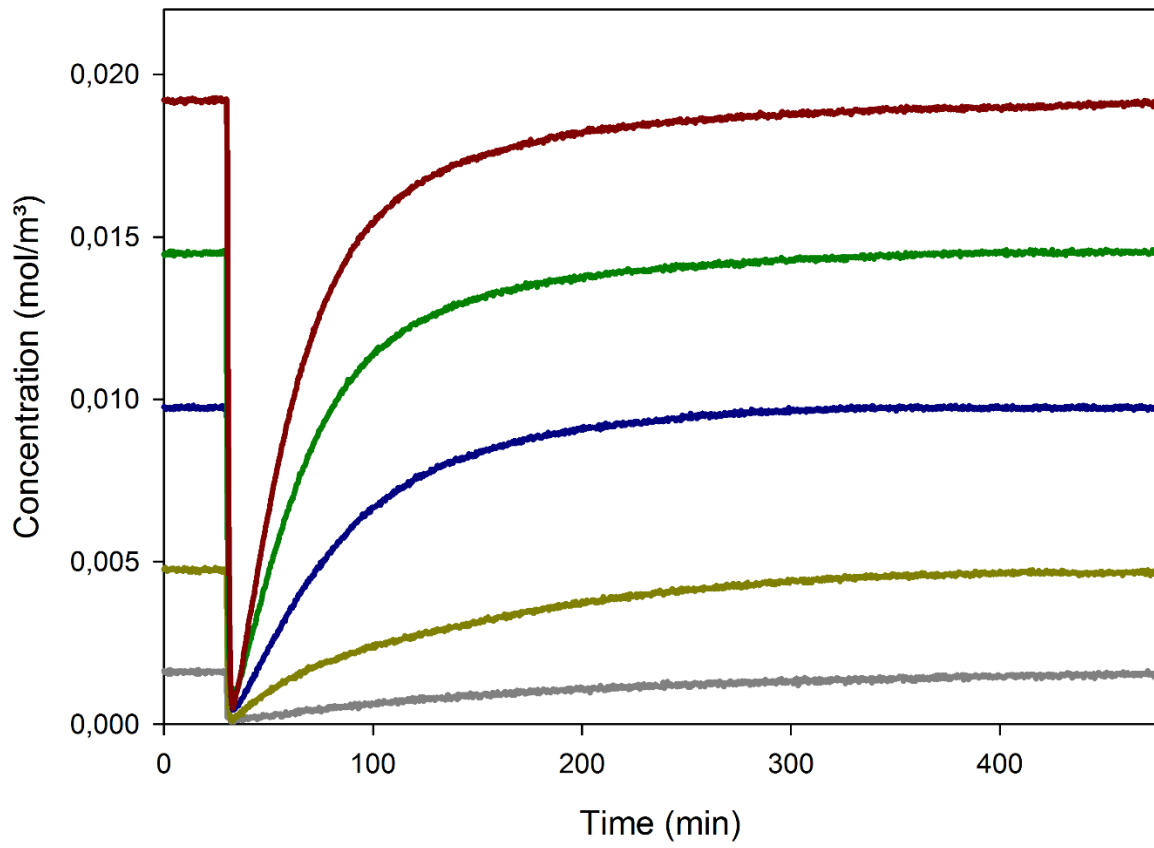
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1 The simulated pressure-velocity curves coincide well with the experimental pressure drop  
2 measurements, as emphasized by coefficients of determination of 0.989, 0.884, 0.955 and 0.817  
3 for 1 to 4 layers of filter respectively. Therefore, Comsol Multiphysics can easily be deployed  
4 to simulate the pressure drop of a filter system at a given flow rate and for a certain system  
5 geometry and filter thickness, with the intrinsic Darcy-Forchheimer parameters.

#### 6 **4.2 Adsorption experiments**

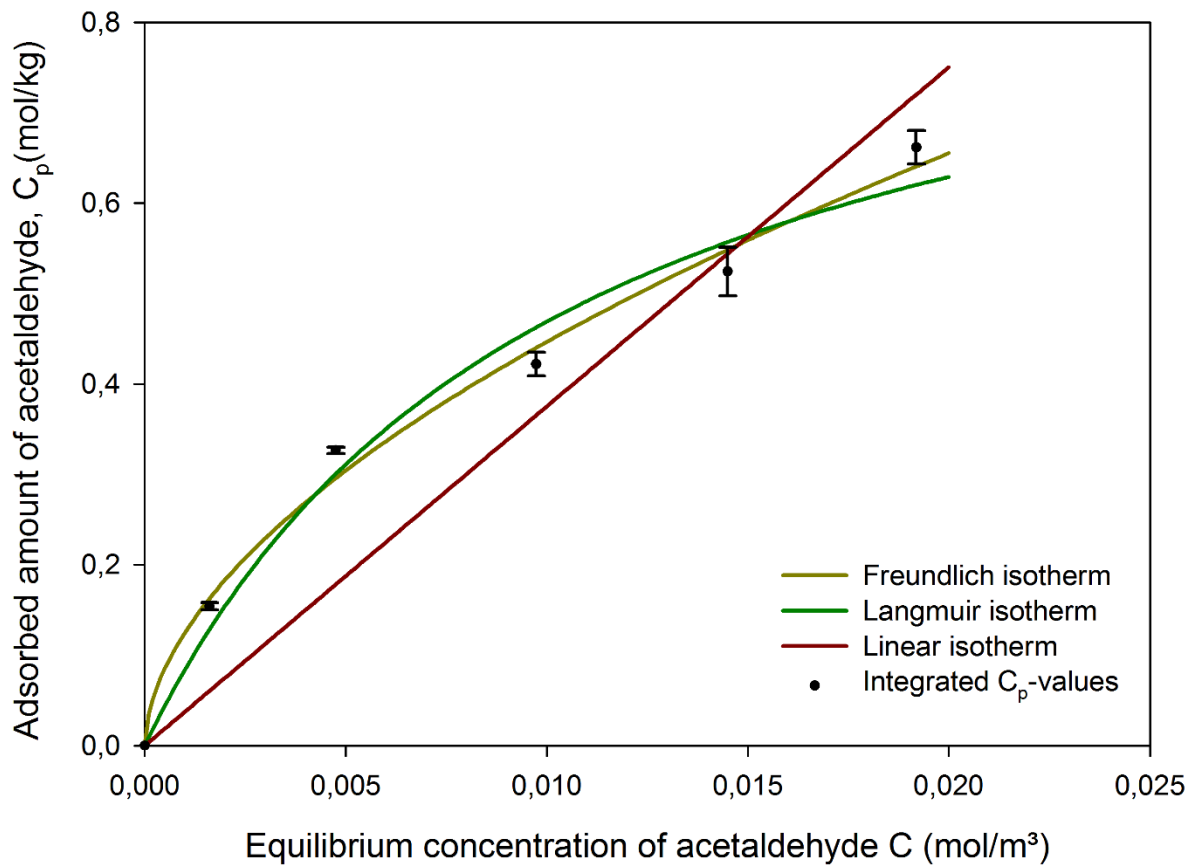
7 The breakthrough curves for 5 different acetaldehyde concentrations were determined (Figure  
8 6). Initial breakthrough is observed almost instantaneously after the start of the experiment, due  
9 to an insufficient layer thickness of the ACF filter and the combined effect of high flow rates,  
10 ppm level concentrations, and the relatively high polarity of acetaldehyde compared to other  
11 indoor air species used in breakthrough experiments with ACF filters, such as aromatic  
12 compounds [15]. However, complete breakthrough lasted until at least 3 hours after initial  
13 breakthrough for all experiments, indicating a large adsorption capacity for acetaldehyde. By  
14 integrating these adsorption curves, the adsorbed amount of acetaldehyde at equilibrium could  
15 be determined. From this, the adsorption isotherm (at  $T=20^{\circ}\text{C}$ ), i.e. the adsorbed amount of  
16 acetaldehyde vs. inlet concentration, could be derived (Figure 7). Although the breakthrough  
17 curves for the repetitive experiments aligned, there was still an uncertainty on the inlet  
18 concentration of acetaldehyde, limited by the mass flow controllers. This effected integration  
19 of the breakthrough curves, as presented by the error bars on Figure 7.

1



**Figure 6** Breakthrough profiles for 5 different acetaldehyde inlet concentrations

2



**Figure 7** Freundlich, Langmuir and linear isotherms fitted with experimentally determined, adsorbed amount of acetaldehyde

1

2 The Langmuir, Freundlich and linear adsorption models could be fitted with the experimental  
 3 integrated  $c_{ads}$ -values with coefficients of determination of 0.980, 0.992 and 0.872 respectively.

4 The resulting parameters are listed in Table 1. Cal et al. (1994) reported values for similar  
 5 pollutants of 0.142-0.580 and 2.81-117 (mol/kg)/(mol/m<sup>3</sup>)<sup>(1/n)</sup> for 1/n and  $K_{Freund.}$  respectively

6 [4]. Yao et al. (2009) fitted the Freundlich equation for 27 different indoor air species, resulting  
 7 in values ranging from 0.072 to 0.821 for 1/n and 0.1 to 7.84 (mol/kg)/(mol/m<sup>3</sup>)<sup>(1/n)</sup> for  $K_{Freund.}$

8 [18]. Brasquet & Le Cloirec (1997) and Popescu et al. (2013) both reported Langmuir  
 9 parameters for several VOCs ranging from 2.12 to 13.84 mol/kg for  $c_{ads,max}$  and 6.44-15000

10 m<sup>3</sup>/mol for  $K_{Langm.}$  [5,15]. Thus, the obtained parameters coincide well with findings in earlier  
 11 studies for similar volatile organic compounds in the same range of pollutant concentrations.

12

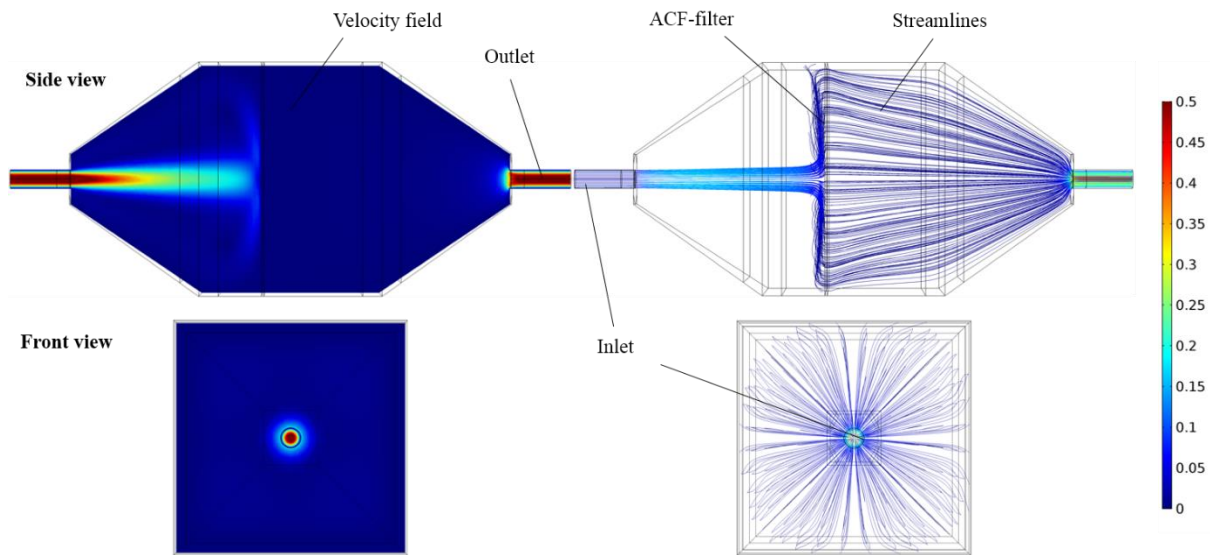


**Table 1** Equilibrium parameters

Parameter	Value
<b>Linear</b>	
$K_{lin.}$ (m <sup>3</sup> /kg)	37.52
<b>Langmuir</b>	
$c_{ads,max}$ (mol/kg)	0.95
$K_{Langm.}$ (m <sup>3</sup> /mol)	96.65
<b>Freundlich</b>	
$1/n$ (-)	0.55
$K_{Freund.}$ (mol/kg)/(mol/m <sup>3</sup> ) <sup>(1/n)</sup>	5.71

### **4.3 Numerical model**

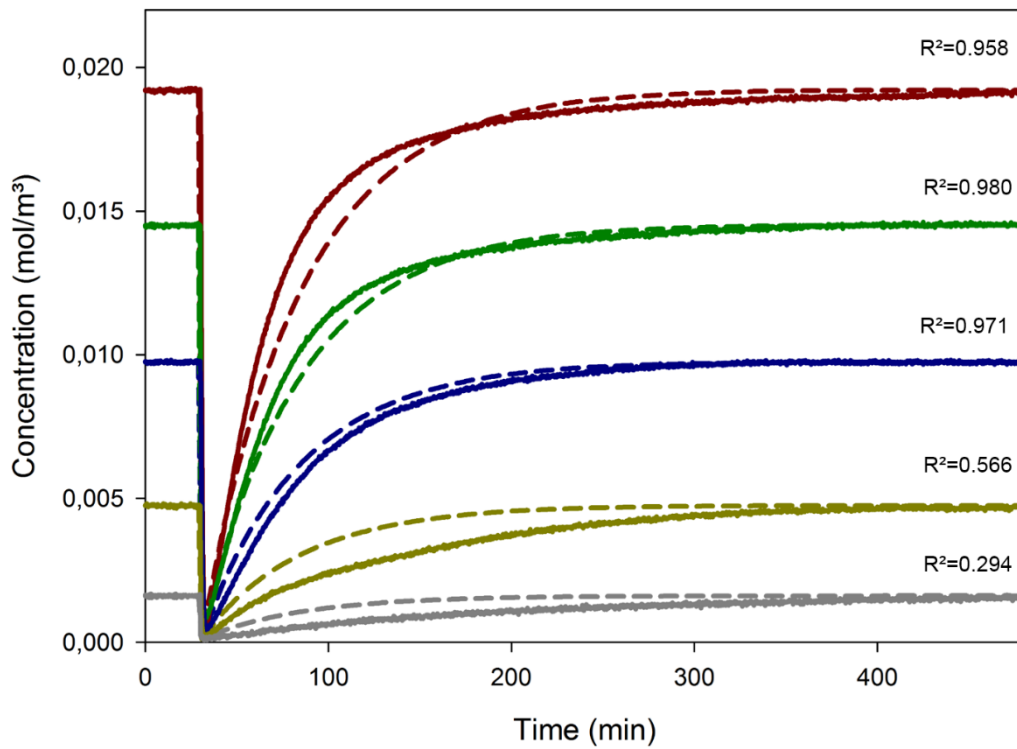
Comsol Multiphysics was also used to simulate the breakthrough curves of acetaldehyde in the filter module. Firstly, the air flow through the filter module was simulated, based on the earlier determined Darcy-Forchheimer parameters, at a flow rate of 200 cm<sup>3</sup>/min. Figure 8 shows that the airflow is uniformly distributed, utilizing the complete filter surface area. However, it should be noted that the high inflow velocity resulted in a concentrated air flow, impacting the filter surface. A more uniform air inflow is desired and should be taken in consideration for upscale purposes. Based on these fluid flow characteristics and the adsorption models with their respective estimated adsorption parameters, the breakthrough curves were simulated for 5 different inlet concentrations and validated against the experimental results.



1

**Figure 8** Fluid flow dynamics (air flow velocity, m/s) inside the ACF-filter module for an airflow rate of 200 cm<sup>3</sup>/min and 0.5 mm filter thickness (1 layer)

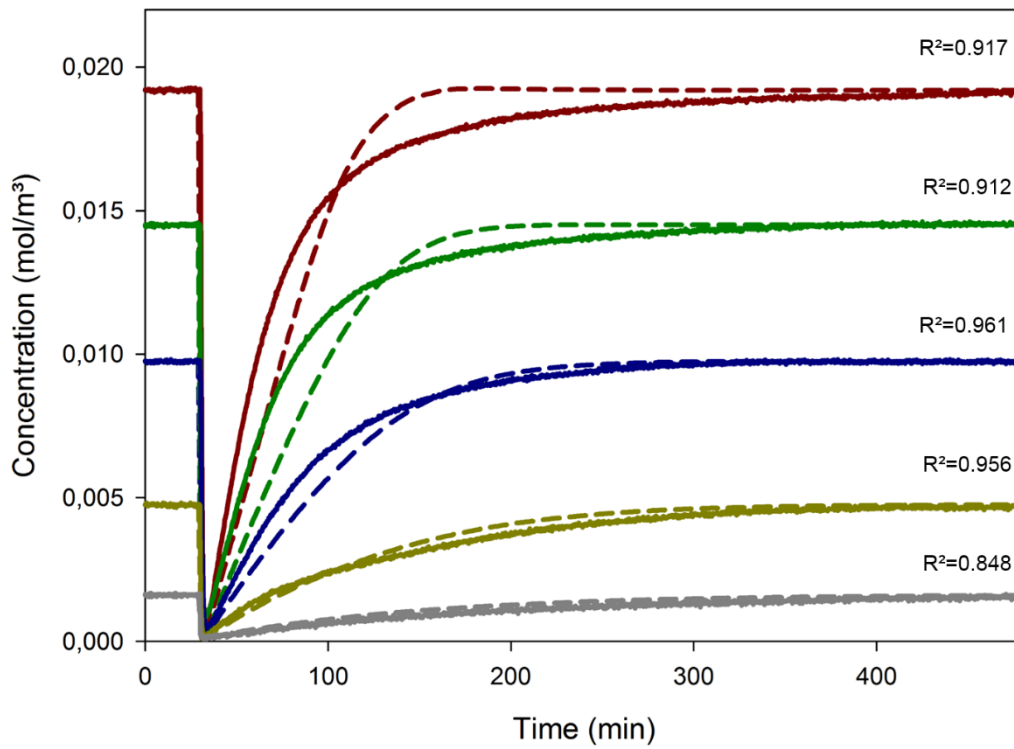
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**Figure 9** Simulated breakthrough profiles with linear isotherm (dashed) vs experimental data (continuous)

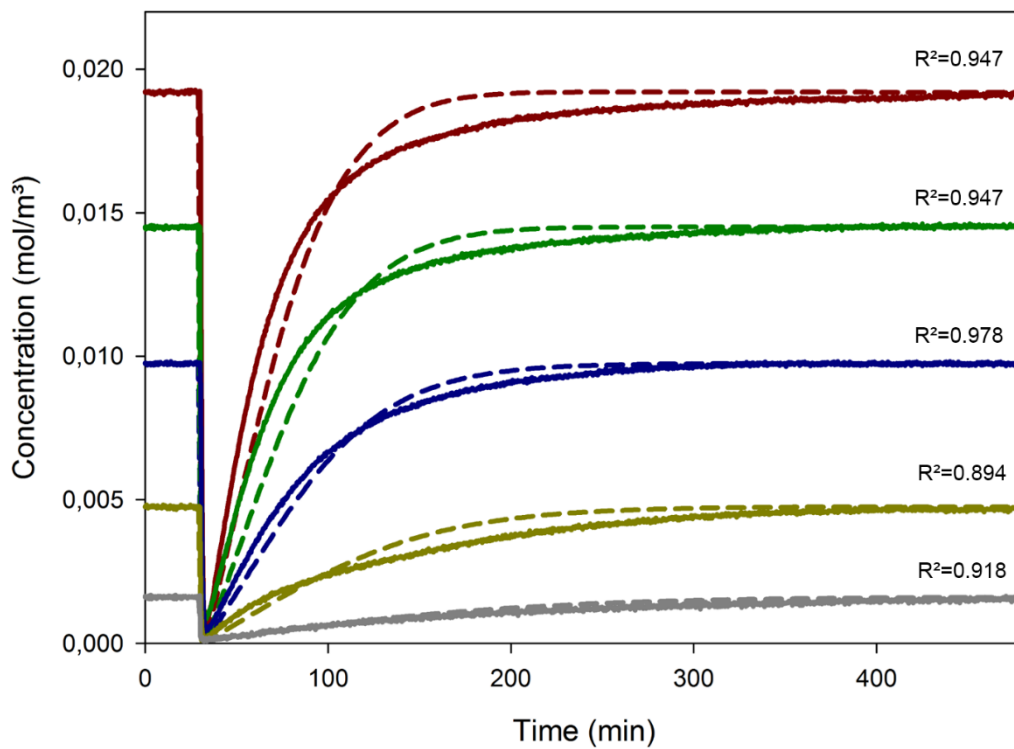
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**Figure 10** Simulated breakthrough profiles with Langmuir isotherm (dashed) vs experimental data (continuous)

2



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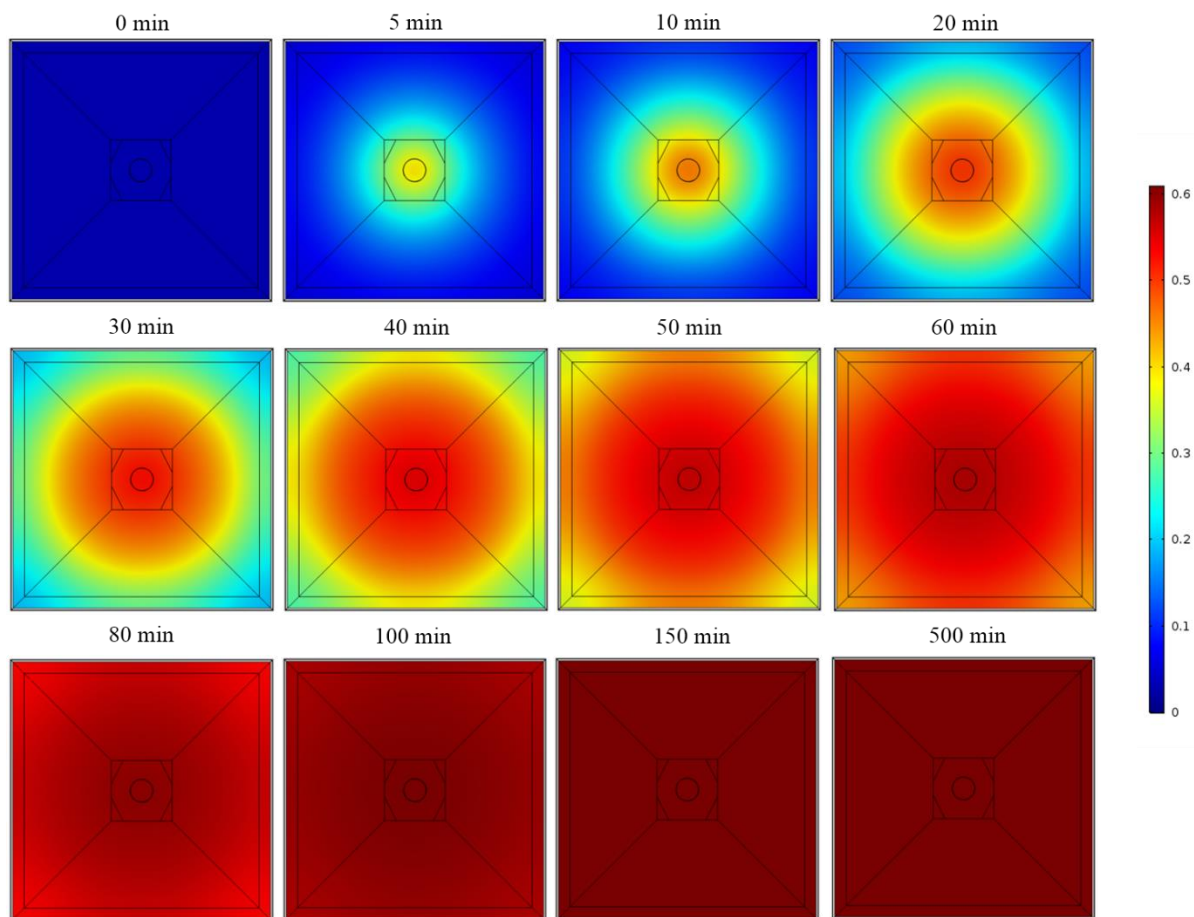
**Figure 11** Simulated breakthrough profiles with Freundlich isotherm (dashed) vs experimental data (continuous)

4

1 Each adsorption model has been assessed by comparing the measured and simulated  
2 breakthrough curves, as shown in Figure 9 to Figure 11. In a first observation, it can be  
3 concluded that the Freundlich adsorption model shows the best overall fit. Secondly, the three  
4 adsorption models deliver significantly different breakthrough curves, as emphasized by the  
5 distinct coefficients of determination. This can be explained by that fact that under dynamic  
6 conditions, the gas phase inside the porous medium may present values that range from 0 to  
7 equilibrium concentration due to mass transfer resistance, resulting in different transient  
8 adsorption curves [22]. Since the model assumes thermodynamic equilibrium, adsorption  
9 occurs instantaneously based on the local concentration and the specific isotherm model.  
10 Including kinetics for adsorption/desorption processes in future work could improve the  
11 simulated breakthrough curves, as we have shown in previous modelling work [21]. This could  
12 clarify the relatively poor fit of Langmuir isotherm simulations for higher inlet concentrations.  
13 The linear model data presents an opposite trend, showing good agreement for higher  
14 acetaldehyde concentrations, but more rapid breakthrough than the experimental data for the  
15 lower inlet concentrations. This could be expected, based on the linear isotherm in Figure 7,  
16 indicating that the linear adsorption model should only be deployed in very small ranges of  
17 equilibrium concentrations. Despite the best fit for the Freundlich model, the simulated data  
18 could not perfectly align with the measured breakthrough curves. The error-level on the  $c_{ads}$ -  
19 values, integrated from the experimental data, will certainly play a role in this and optimization  
20 of the adsorption parameters could further improve the model. However, the main reason for  
21 the deviation of the simulated breakthrough curve is possibly due to the influence of  
22 intraparticle resistance on the mass transfer kinetics. Several studies show that intraparticle  
23 mass transfer is the limiting step of the adsorption kinetics, while distinguishing between pore  
24 diffusion, as defined by the effective diffusion coefficient, and surface diffusion [16,17,22,23].  
25 The latter becomes increasingly important with a higher microporous volume, as can be the

1 case for activated carbon fiber filters. Unfortunately, Comsol Multiphysics v5.3 does not  
2 provide an option to include surface diffusion in the governing equations, so that mass transfer  
3 inside the pores can only be expressed by pore volume diffusion. Moreover, it could be  
4 interesting, to assess alternative adsorption isotherms in future research. For example the  
5 empirical Dubinin-Radushkevich isotherm model, which is used to describe adsorption of  
6 subcritical vapors onto microporous solids such as activated carbon or zeolites [24,25].  
7 Additionally, better fittings with experimental data could be achieved by isotherm models  
8 which combine both Langmuir and Freundlich expressions, such as Sips or Koble-Corrigan  
9 isotherm models [24] or the Toth model, which is an improvement to the Langmuir isotherm  
10 model [24,26,27]. Nevertheless, the developed model has proven acceptable performance,  
11 given the complex dynamic behavior of the system, and it can provide valuable insights such  
12 as spatial variation in saturation level of the level, as shown in Figure 12.

13



**Figure 12** Spatial variation of adsorbed acetaldehyde (mol/kg) on the ACF filter (as seen from the inlet) at different time intervals (min); simulated with the Freundlich model.

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2  
3 It can be seen that the dynamics of the fluid flow greatly influences the spatial saturation of the  
4 ACF filter. Saturation starts in the center of the filter, due to the high inlet velocity, and spreads  
5 towards the edges of the filter later in the adsorption process. The saturation process can directly  
6 be linked to the breakthrough curves in Figures 9 to 11. The steep drop in the breakthrough  
7 curves results from a rapid adsorption in the center of the filter, followed by rapid increase in  
8 the outlet concentration of the filter module. After 1 hour, the curves begin to flatten out towards  
9 equilibrium concentration. This is reflected by the mere change in saturation level after 1 hour  
10 as seen in Figure 12, i.e. a color change from light red to dark red. Based on the spatial variation  
11 of adsorbed acetaldehyde combined with fluid flow dynamics, the geometry of the filter  
12 module can be adapted and optimized in the numerical model to achieve a more uniform level

1 of saturation during the operating time of the filter and hence, a better control over the  
2 breakthrough curves.

### 3 **5. Conclusions**

4 The pressure drop over an ACF-filter can be described with the Darcy-Forchheimer law with  
5 unique parameter values for various filter thicknesses. The parameters were determined by  
6 quadratic regression, using pressure-velocity measurements in a closed circuit vent duct. CFD-  
7 modelling has successfully demonstrated the validity of the Darcy-Forchheimer parameters.  
8 Consequently, the CFD model was extended to simulate the breakthrough profile of  
9 acetaldehyde adsorption on the ACF-filter. Three adsorption isotherm models were  
10 investigated: Linear, Langmuir and Freundlich adsorption. The Freundlich model has proven  
11 the most suitable fit with experimental breakthrough profiles for different inlet concentrations  
12 of acetaldehyde. However, there is still some room for improvement by including surface  
13 diffusion as a mass transfer mechanism inside the pore system, combined with effective pore  
14 diffusion. In this work, only the latter could be included due to practical limitations.  
15 Nonetheless, the numerical model was able to generate a good agreement with experimental  
16 data and could provide valuable insights. Continuous improvement of the numerical model will  
17 speed up the development of commercial ACF devices for indoor air purification.

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## 1 6. Appendix

$c$	<i>Concentration (mol/m<sup>3</sup>)</i>
$C_{ads}$	<i>Amount of adsorbed acetaldehyde (mol/kg)</i>
$C_{ads,max}$	<i>Maximum amount amount of adsorbed acetaldehyde (mol/kg)</i>
$D_e$	<i>Effective diffusivity (m<sup>2</sup>/s)</i>
$D_K$	<i>Knudson diffusion coefficient [m<sup>2</sup>/s]</i>
$I$	<i>Identity matrix</i>
$K_{Freund.}$	<i>Freundlich equilibrium constant (mol/kg)/(m<sup>3</sup>/mol)<sup>n</sup></i>
$K_{Langm.}$	<i>Langmuir equilibrium constant (m<sup>3</sup>/mol)</i>
$K_{lin.}$	<i>Linear equilibrium constant (m<sup>3</sup>/kg)</i>
$1/n$	<i>Freundlich exponent (–)</i>
$P$	<i>Pressure (Pa)</i>
$v$	<i>Velocity (m/s)</i>
$x$	<i>Filter thickness (m)</i>
$\beta$	<i>Forchheimer term (1/m)</i>
$\epsilon_p$	<i>Porosity (–)</i>
$\kappa$	<i>Permeability (m<sup>2</sup>)</i>
$\mu$	<i>Dynamic viscosity [Pa · s]</i>
$\mu_T$	<i>Eddy viscosity [Pa · s]</i>
$\rho$	<i>Density [kg/m<sup>3</sup>]</i>
$\tau$	<i>Tortuosity (–)</i>

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## 7. References

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